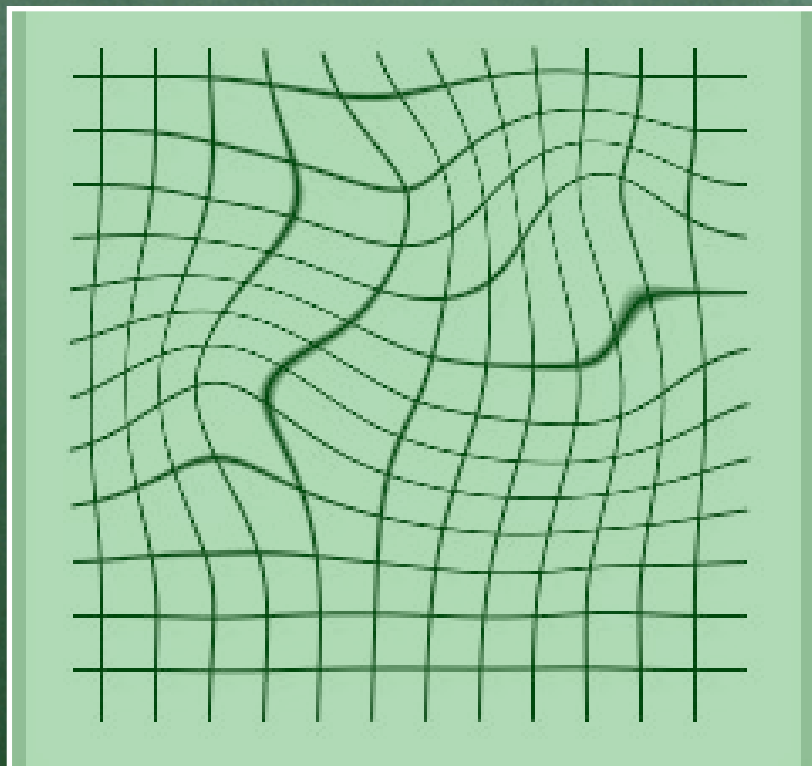


# Nonperturbative quantum field theory on a real time lattice

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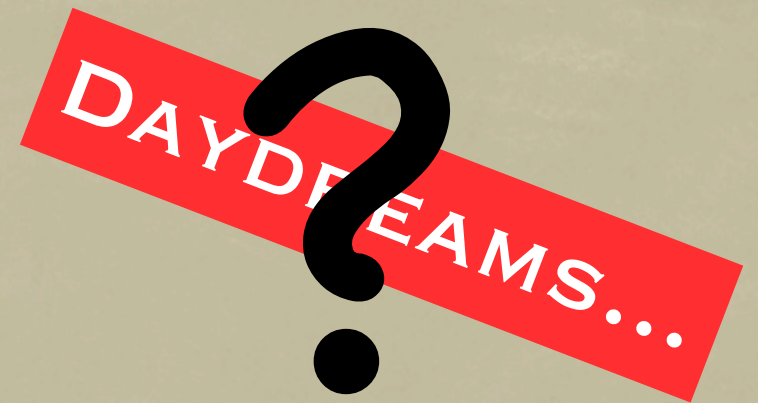
# *Real time simulations?*

DAYDREAMS...

- spectral function, decay rates
- real time response  $\rightarrow$  nonequilibrium
- nonequilibrium field theory without
  - Rayleigh-Jeans problem (classical)
  - Gauge dependence problems
- could we simulate a heavy ion collision?  
*(short period of time: a few fm/c would be enough)*



# *Real time simulations?*



- spectral function, decay rates
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*(short period of time: a few fm/c would be enough)*

$\rightarrow$  stochastic quantization techniques



# Stochastic quantization

Parisi, Wu 1981

$$\partial_{\vartheta} \phi(x, \vartheta) = -\frac{\delta S_E[\phi]}{\delta \phi(x, \vartheta)} + \eta(x, \vartheta)$$

$$\langle \eta(x_1, \vartheta_1) \eta(x_2, \vartheta_2) \rangle = 2\delta(\vartheta_1 - \vartheta_2) \delta^{(4)}(x_1 - x_2)$$

$$\phi(\vartheta + \varepsilon) = \phi(\vartheta) - \varepsilon \frac{\partial S_E}{\partial \phi} + \sqrt{2\varepsilon} \xi$$

$$\langle \xi^2 \rangle = 1 \text{ random number}$$

$$\frac{\partial \Phi(\mathbf{x}, t)}{\partial t} = i \frac{\delta S[\Phi]}{\delta \Phi(\mathbf{x}, t)} + \eta(\mathbf{x}, t)$$

Huffel, Rumpf 1984

Gozzi 1984

Klauder 1984

$$\varphi'_R(t, \mathbf{x}) = \varphi_R(t, \mathbf{x}) - \epsilon I_\xi(\varphi_R, \varphi_I; t, \mathbf{x}) + \sqrt{\epsilon} \eta_R(t, \mathbf{x})$$

$$\varphi'_I(t, \mathbf{x}) = \varphi_I(t, \mathbf{x}) + \epsilon R_\xi(\varphi_R, \varphi_I; t, \mathbf{x}) + \sqrt{\epsilon} \eta_I(t, \mathbf{x})$$

$$R_\xi(\varphi_R, \varphi_I; t, \mathbf{x}) \equiv \text{Re} \left( \frac{\delta S_\xi[\varphi]}{\delta \varphi(t, \mathbf{x})} \Big|_{\varphi=\varphi_R+i\varphi_I} \right)$$

$$I_\xi(\varphi_R, \varphi_I; t, \mathbf{x}) \equiv \text{Im} \left( \frac{\delta S_\xi[\varphi]}{\delta \varphi(t, \mathbf{x})} \Big|_{\varphi=\varphi_R+i\varphi_I} \right)$$

Real Fokker-Planck equation

$$\frac{\partial P[\varphi_R, \varphi_I]}{\partial \vartheta} = \int d^4x \left[ \frac{\delta(PI)}{\delta \varphi_R} - \frac{\delta(PR)}{\delta \varphi_I} + \frac{\delta^2 P}{\delta \varphi_R^2} \right]$$

# Reweighting vs. Stochastic quantization

## Reweighting

$$\int D\Phi e^{iS_R[\Phi]-S_I[\Phi]} \Phi(t_1)\Phi(t_2) = \int D\Phi e^{-S_I[\phi]} \left[ e^{iS_R[\Phi]} \Phi(t_1)\Phi(t_2) \right]$$

Action in Minkowski time: no importance sampling.

Cost of simulation:  $\sim \exp(\text{volume})$

## Stochastic quantization *(in real time)*

Analytical continuation of the distribution!

Strict argument for convergence in free theory.

Thermalization time? Convergence? Precision test?

Performance compared to stochastic quantization

Callaway et al. (1985)



# The real field becomes complex

$$\langle \mathcal{O} \rangle_{\vartheta} = \int [d\varphi_R][d\varphi_I] \mathcal{O}(\varphi_R + i\varphi_I) P(\varphi_R, \varphi_I, \vartheta)$$

$$\langle \mathcal{O} \rangle_{\vartheta} = \int [d\varphi_R][d\varphi_I] \mathcal{O}(\varphi_R) P_{\text{eff}}(\varphi_R, \vartheta)$$

$$P_{\text{eff}}(\varphi_R, \vartheta) = \int [d\varphi_I] P(\varphi_R - i\varphi_I, \varphi_I, \vartheta)$$

$$P_{\text{eff}}(\varphi_R, \vartheta) \rightarrow e^{iS[\varphi_R]}$$

Lattice in Minkowski  
space-time:

needs a small regulator  $m^2 - i\epsilon$

boudary conditions?

initial conditions?



Real time observables:

$$\text{Tr} \hat{\rho} \hat{\phi} e^{i\hat{H}t} \hat{\phi} e^{-i\hat{H}t}$$

Real time equilibrium:

$$\text{Tr} e^{-\beta \hat{H}} \hat{\phi} e^{iHt} \hat{\phi} e^{-iHt}$$

Contour?

# CTP → Complex Time Path

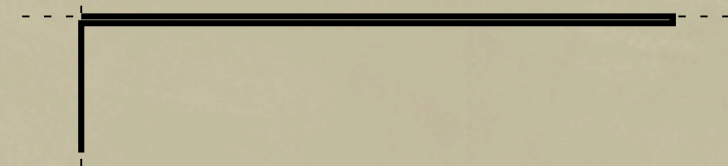
$$-iS = -\frac{i}{2} \sum_j \frac{(\phi_{j+1} - \phi_j)^2}{C_{j+1} - C_j} + \frac{i}{2} \sum_j (C_{j+1} - C_j) (V[\phi(C_{j+1})] + V[\phi(C_j)])$$

Langevin equation:

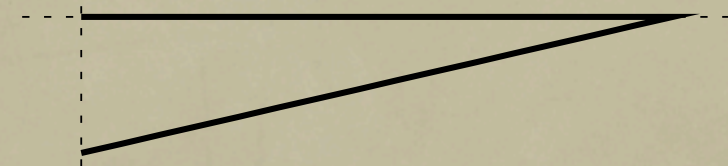
$$\frac{\partial \phi(C_j)}{\partial \vartheta} = i \frac{\partial S}{\partial \phi(C_j)} + \eta_j(\vartheta)$$

Contour points:  $C_j$

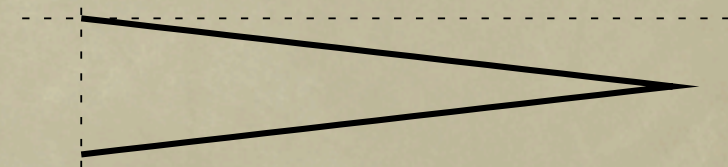
real-time



right

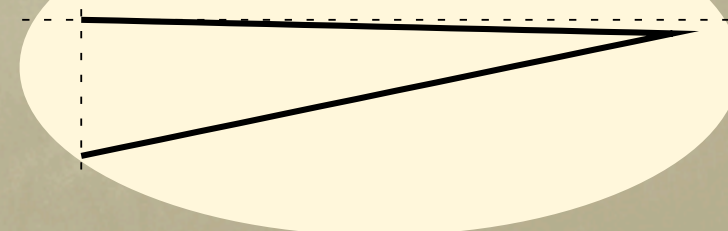


equilateral

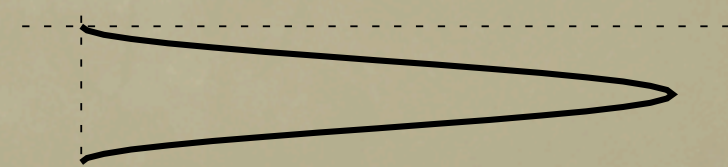


Real time

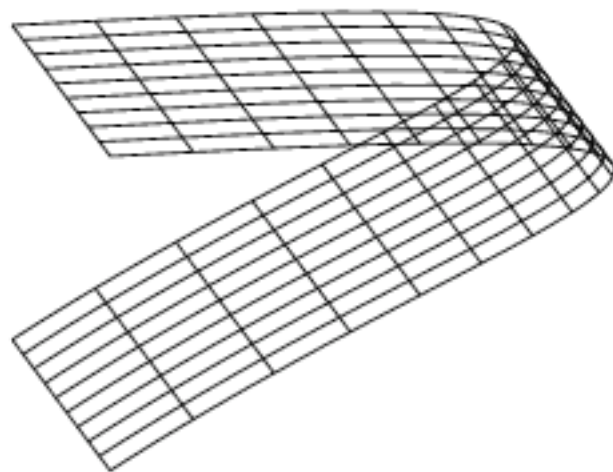
asymmetric



cosine



***Respects gauge symmetry!***



Euclidean time





# Harmonic Oscillator:

Klauder 1984, Nakazato&Yamanaka 1985

$$-iS[\phi] = \frac{1}{2} \phi_i M_{ij} \phi_j$$

$$\phi_j \equiv \phi(C_j)$$

$M$ : complex symmetric matrix

$$\partial_{\vartheta} \phi_i = -M_{ij} \phi_j + \eta_i$$

$$\phi_i = \sum_a z_a \psi_a^i$$

Eigensystem:  $\vec{\psi}_a$  (orthogonal), eigenvalues  $\lambda_a$  (complex)

$$\partial_{\vartheta} z_a = -\lambda_a z_a + \eta'_a$$

Converges if  $\text{Re } \lambda > 0$

$$\text{Re } \langle z^2 \rangle \rightarrow \text{Re } 1/\lambda$$

$$\text{Im } \langle z^2 \rangle \rightarrow \text{Im } 1/\lambda$$

$$\langle z^2 \rangle - 1/\lambda \sim e^{-2\lambda t}$$

$$\langle z^4 \rangle_0 = 3\langle z^2 \rangle_0$$

Complex eigenvalue:

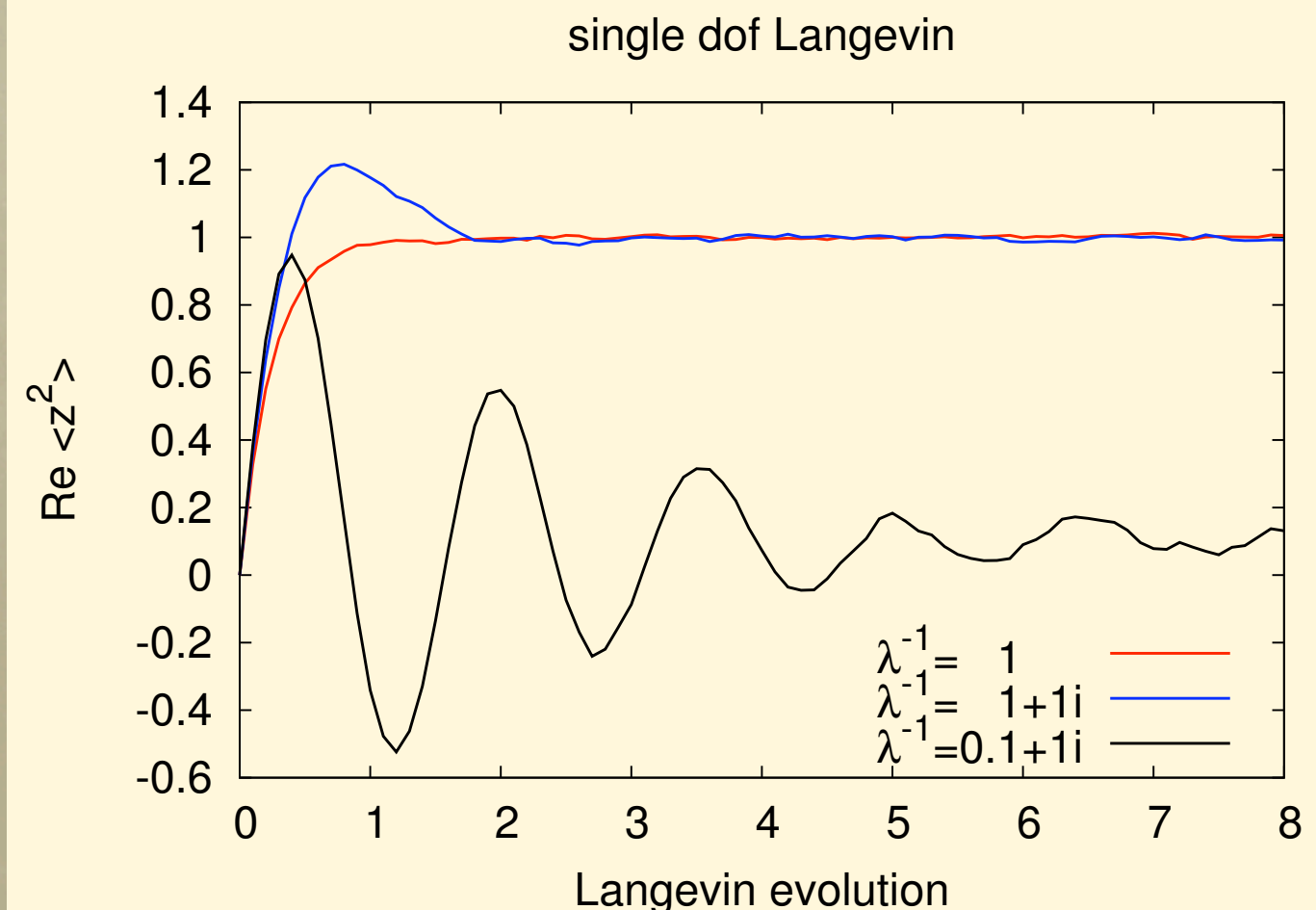
damped oscillation around the limit

Field theory: summing over modes

“Prethermalization”:

accelerates convergence

*Eigenvalues depend on the contour shape.  
The tilt of the contour acts as regulator.*





# Free field: analytically accessible

Nakazato&Yamanaka 1985

$$\dot{\phi}_R(k,t) = -q\phi_I(k,t) - \varepsilon\phi_R(k,t) + \eta(k,t),$$

$$\dot{\phi}_I(k,t) = q\phi_R(k,t) - \varepsilon\phi_I(k,t),$$

$\varepsilon$  small real part  
of the eigenvalue  $\lambda_k = \varepsilon - iq$

## FP Hamiltonian

$$\frac{\partial}{\partial t}P[\phi;t] = H[\phi]P[\phi;t],$$

$$H[\phi] = \int d^4k \left\{ \frac{\delta}{\delta\phi_R(k)} \frac{\delta}{\delta\phi_R(-k)} + \frac{\delta}{\delta\phi_R(k)} [q\phi_I(k) + \varepsilon\phi_R(k)] + \frac{\delta}{\delta\phi_I(k)} [-q\phi_R(k) + \varepsilon\phi_I(k)] \right\}$$

$$\begin{aligned} P_{eq}[\phi] &= \lim_{t \rightarrow \infty} P[\phi;t] \\ &= N' \exp[-\varepsilon \int d^4k \{ |\phi_R(k)|^2 + [1 + \varepsilon^2/(\varepsilon^2 + q^2)] |\phi_I(k)|^2 \\ &\quad - 2(\varepsilon/q)\phi_R(k)\phi_I(-k) \} ] . \end{aligned}$$

$$P_{eff}(\varphi_R, \vartheta) = \int [d\varphi_I] P(\varphi_R - i\varphi_I, \varphi_I, \vartheta)$$

$$\langle \phi(k)\phi(k') \rangle$$

$$= \langle \phi_R(k)\phi_R(k') - \phi_I(k)\phi_I(k') \rangle + i \langle \phi_R(k)\phi_I(k') + \phi_I(k)\phi_R(k') \rangle$$

$$= \delta^4(k+k') \left( \frac{\varepsilon}{q^2 + \varepsilon^2} + i \frac{q}{q^2 + \varepsilon^2} \right) = \delta^4(k+k') \frac{i}{q + i\varepsilon} = \delta^4(k+k') \frac{i}{k^2 - m^2 + i\varepsilon}$$



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## FP Hamiltonian

$$H[\phi] = \int d^4k \left\{ \frac{\delta}{\delta\phi_R(k)} \frac{\delta}{\delta\phi_R(-k)} + \frac{\delta}{\delta\phi_R(k)} [q\phi_I(k) + \varepsilon\phi_R(k)] + \frac{\delta}{\delta\phi_I(k)} [-q\phi_R(k) + \varepsilon\phi_I(k)] \right\}$$

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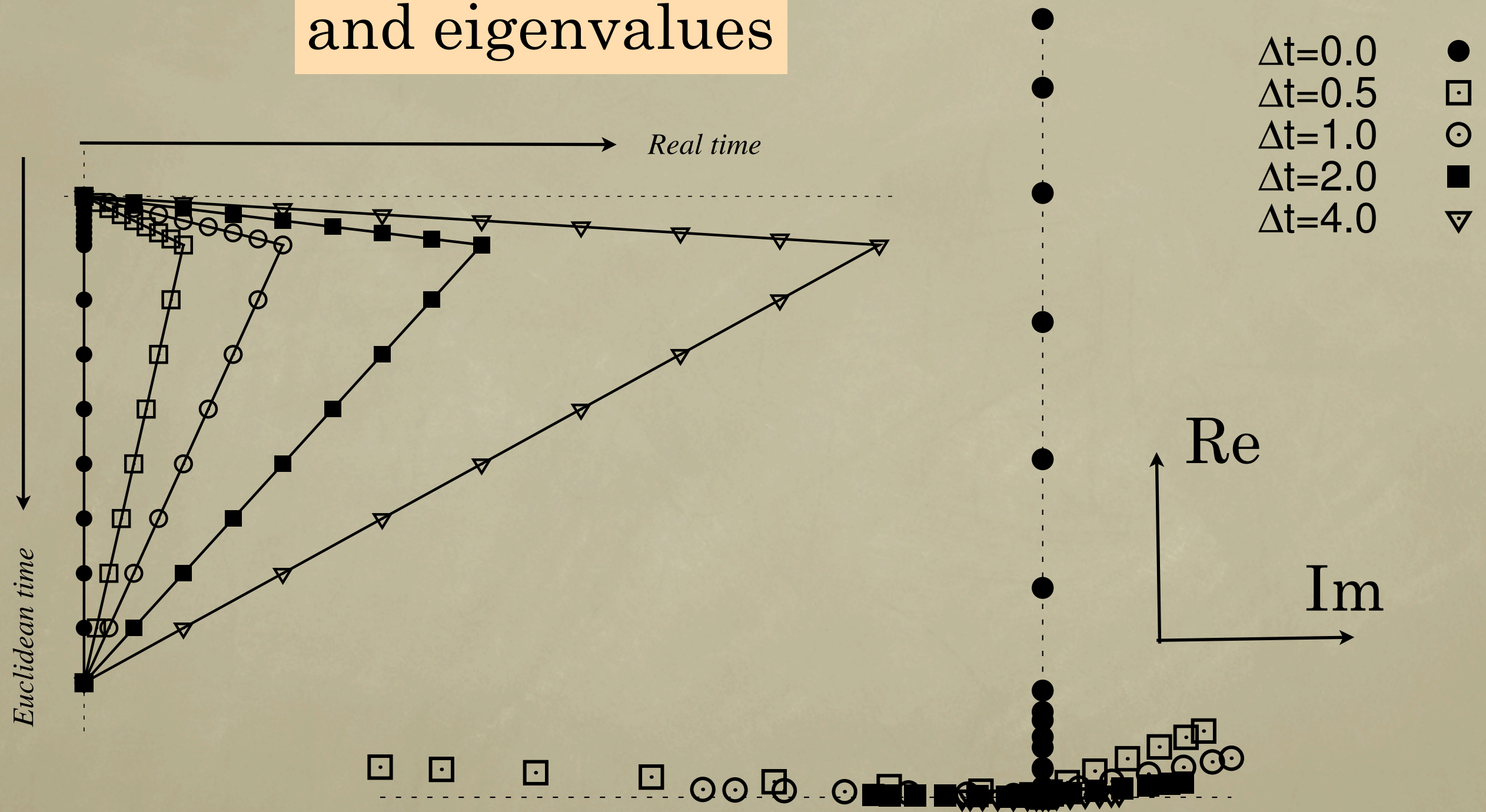
Perturbative propagator:

$\psi_x^a$ : eigenfunctions of  $G_0^{-1}(x, y)$

$$\langle \phi(x)\phi(y) \rangle_0 = \langle z^a z^b \rangle_0 \psi_x^a \psi_y^b = \frac{\delta^{ab}}{\lambda^a} \psi_x^a \psi_y^b = \sum_a \frac{1}{\lambda^a} \psi_x^a \psi_y^a = [G_0^{-1}(x, y)]^{-1}$$



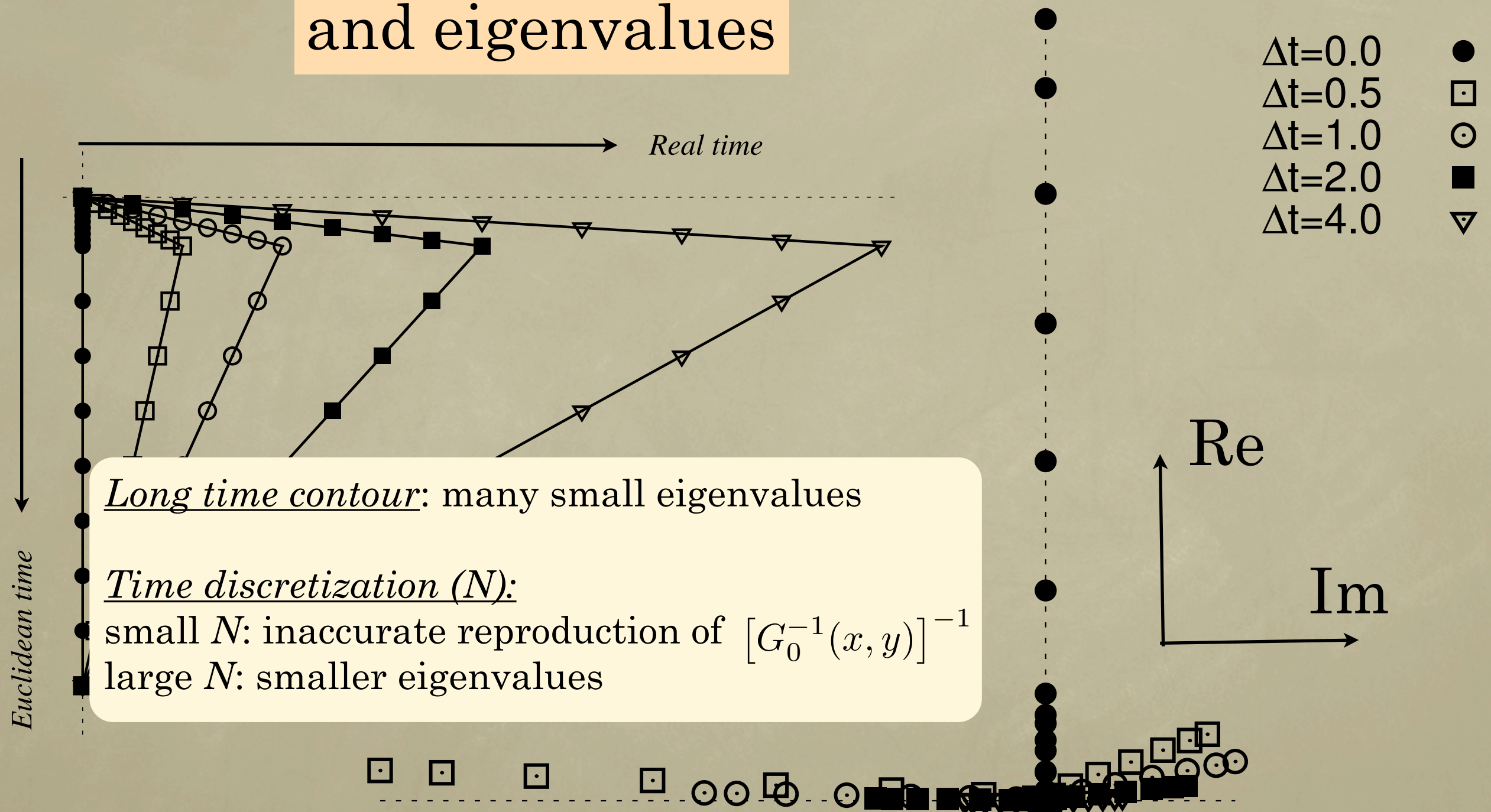
# Contour shapes and eigenvalues



The real part of the eigenvalues are always positive if there is a tilt, no matter how small

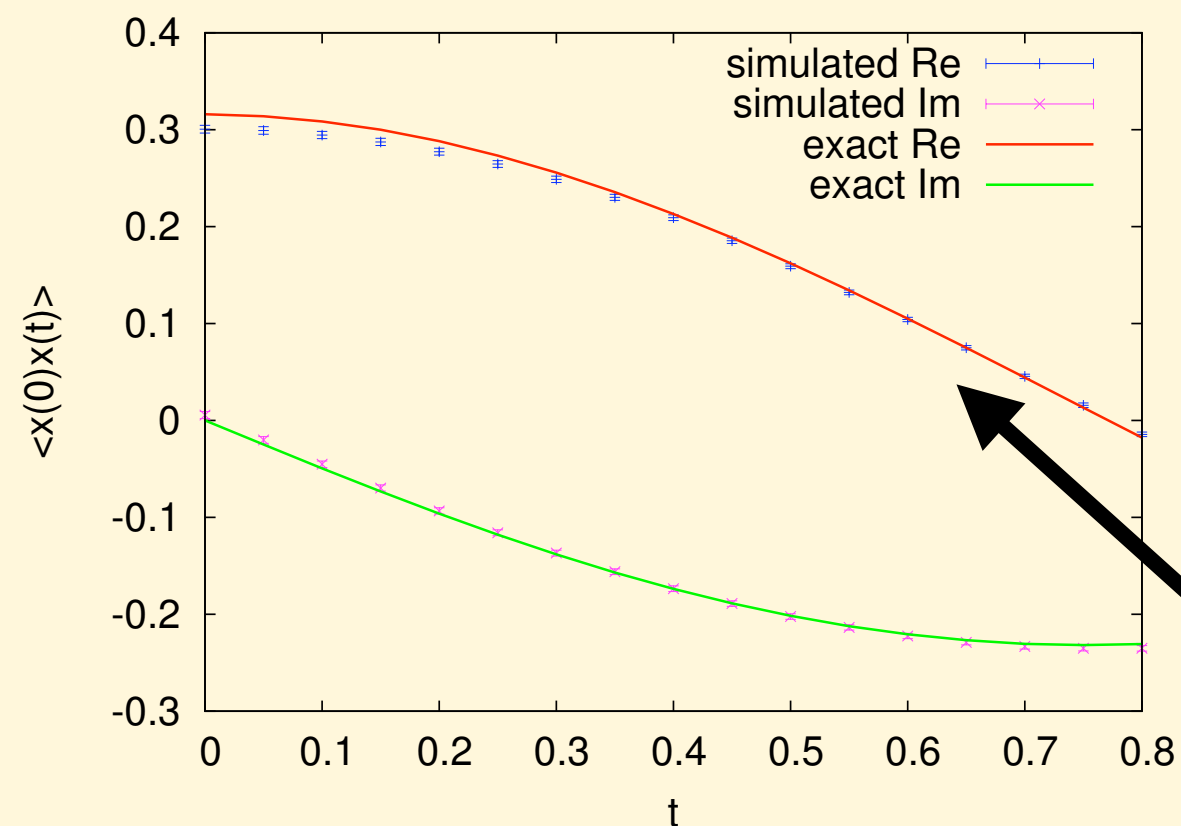
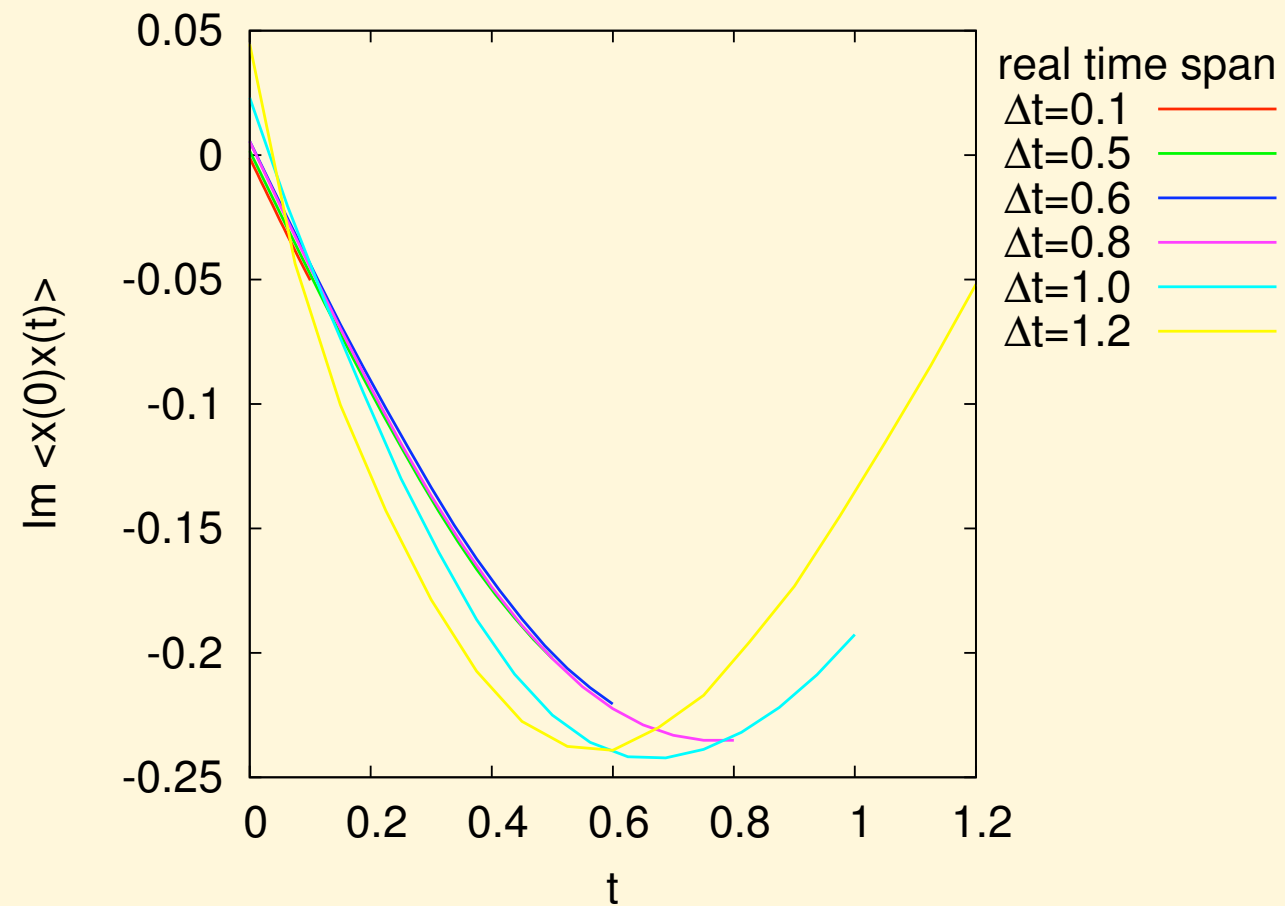


# Contour shapes and eigenvalues



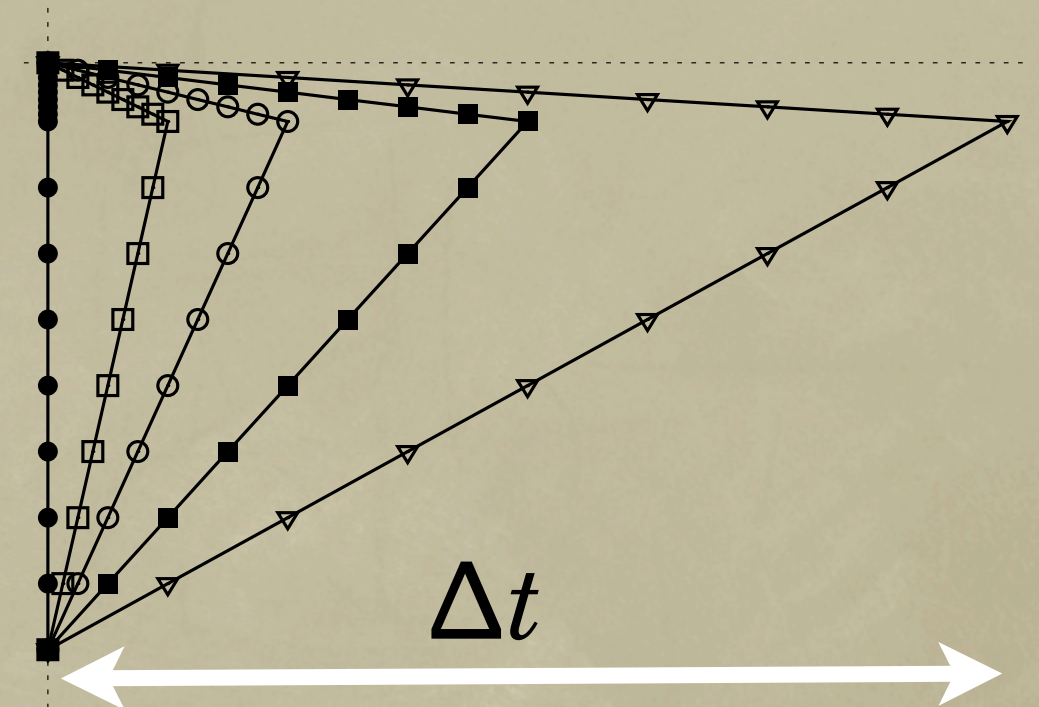
The real part of the eigenvalues are always positive if there is a tilt, no matter how small





## Short time intervals

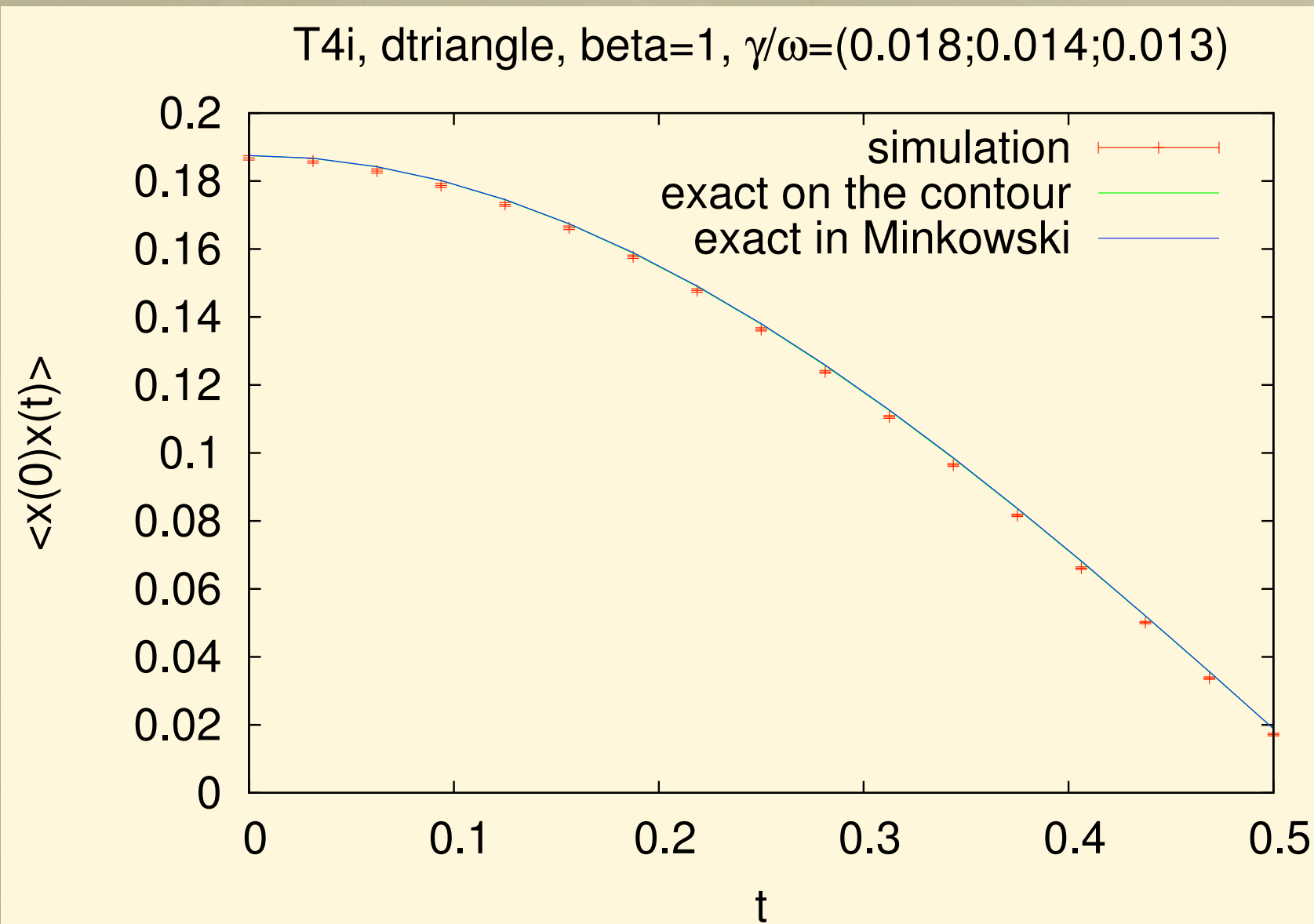
$$\langle \phi(0)\phi(t) \rangle \quad t = 0 \dots \Delta t$$



Results independent of real-time span of the grid until a certain limit

Schrodinger's exact result on the contour

# A precision test: Extracting damping rate

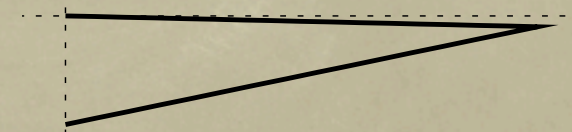


Not exactly real time:  
tilt angle=0.002

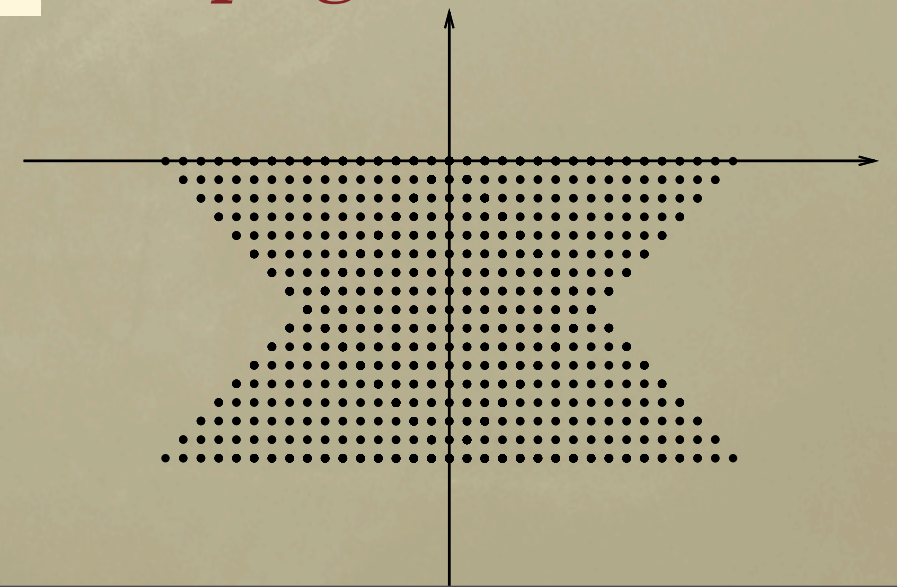
$\gamma/\omega < 2\%$       accuracy 40%

Fitting  $\sim \cos(\omega t)e^{-\gamma t}$

Requires precision data



*Propagator known at:*

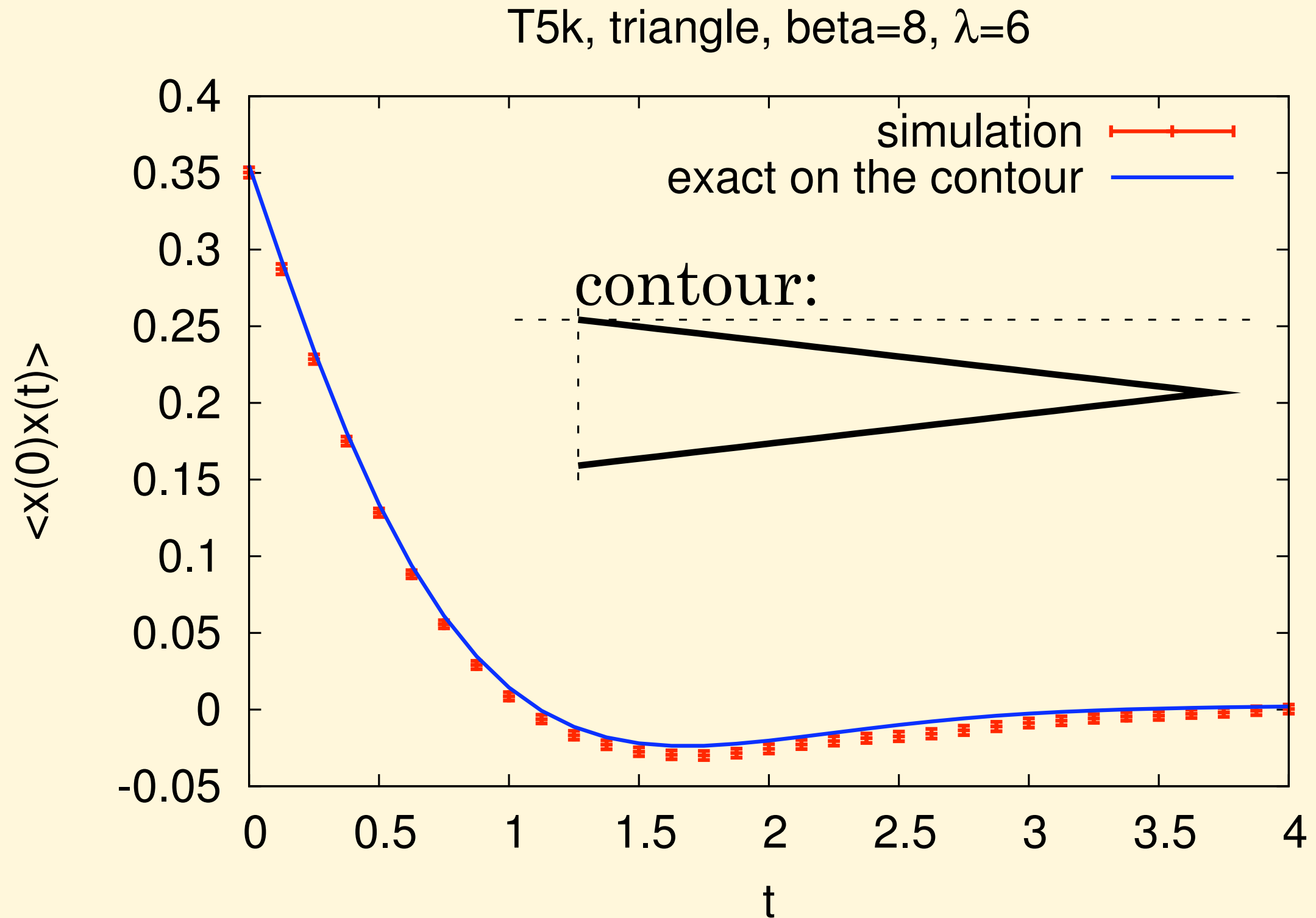


Possible improvements:

statistical analysis based on the  
correlation matrix

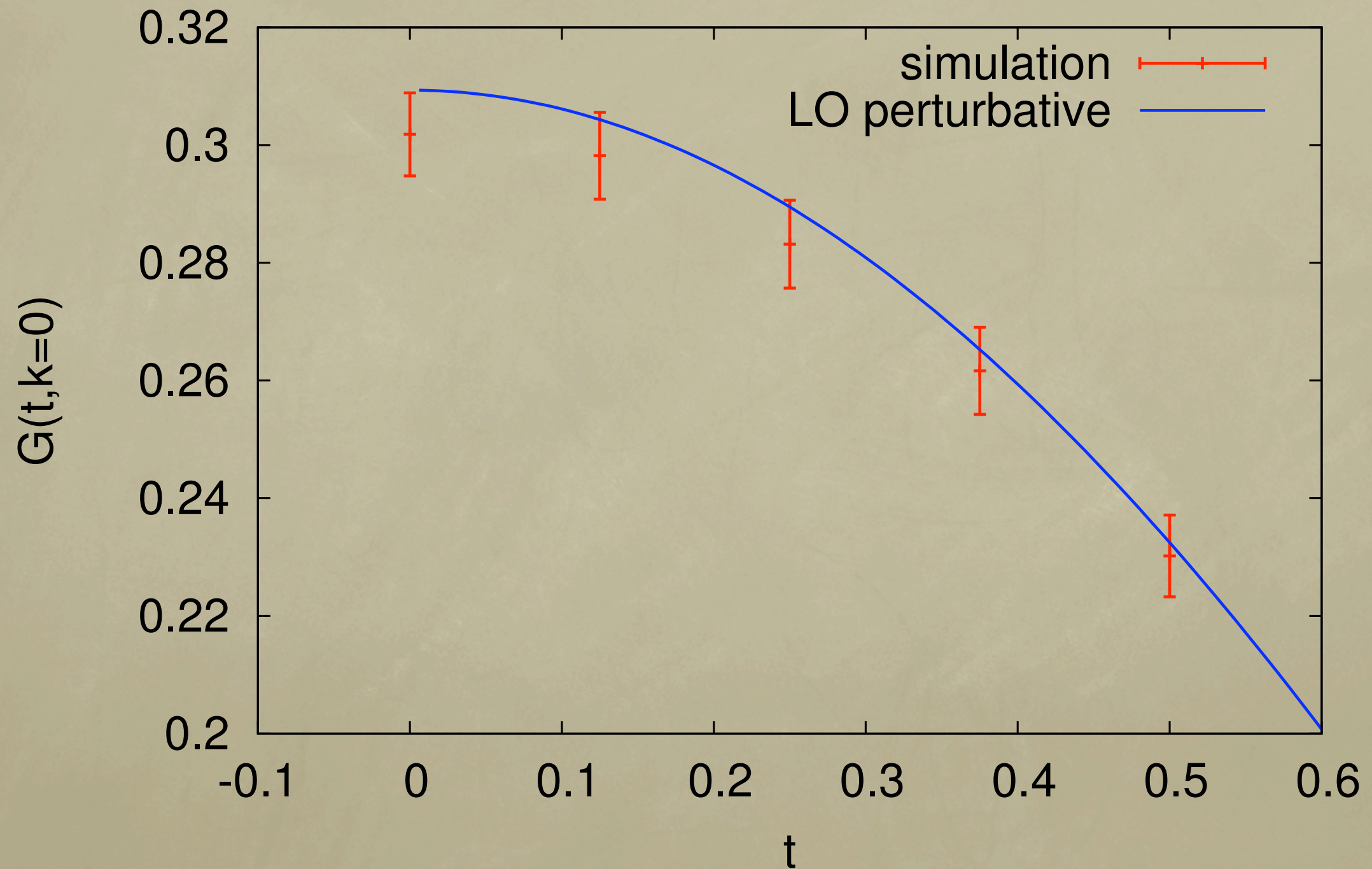


At lower temperature: longer intervals



# Field theory:

1+1 dim. scalar field theory

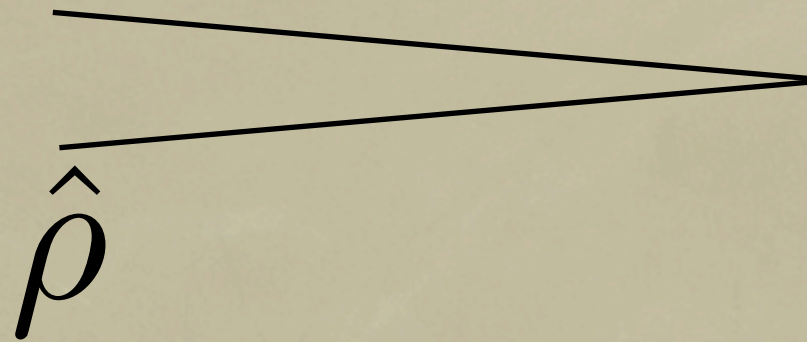


*solid curve: Hartree approx.*



Nonequilibrium:  $\text{Tr} \hat{\rho} \hat{\phi} e^{i\hat{H}t} \hat{\phi} e^{-i\hat{H}t}$

The contour:



*Tilt: regulator.  
(modified initial  
condition and  
more damping)*

Simplest: *Gaussian density operator*

$$S'[\phi_+, \phi_-] = S[\phi_+] - S[\phi_-] + \frac{1}{a_t} S_0(\phi_+(t_i), \phi_-(t_i))$$

$$\begin{aligned} S_0[\phi_+, \phi_-] &= i\dot{\Phi}(\phi_+ - \phi_-) - \frac{\sigma^2 + 1}{8\xi^2} \left( (\phi_+ - \Phi)^2 + (\phi_- - \Phi)^2 \right) \\ &+ \frac{i\eta}{2\xi} \left( (\phi_+ - \Phi)^2 - (\phi_- - \Phi)^2 \right) \\ &+ \frac{\sigma^2 - 1}{4\xi^2} (\phi_+ - \Phi)(\phi_- - \Phi) \end{aligned}$$

Field theory:

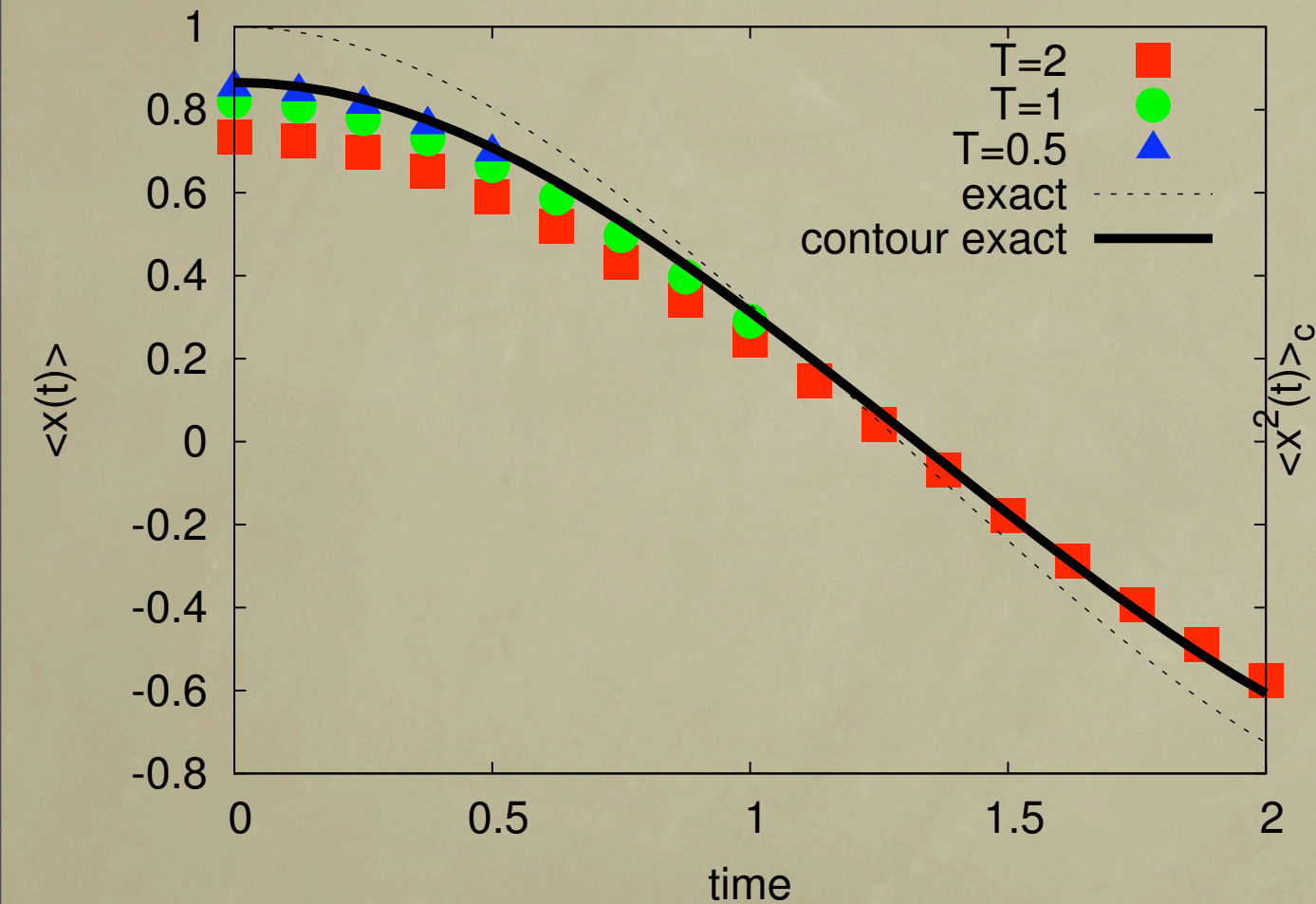
Fourier transformation in each  
Langevin time step

$$\begin{aligned} \Phi &= \langle \phi(t_i) \rangle \\ \dot{\Phi} &= \langle \dot{\phi}(t_i) \rangle \\ \xi^2 &= \langle \phi(t_i) \phi(t_i) \rangle_c \\ \eta\xi &= \frac{1}{2} \langle \dot{\phi}(t_i) \phi(t_i) + \phi(t_i) \dot{\phi}(t_i) \rangle_c \\ \eta^2 + \frac{\sigma^2}{4\xi^2} &= \langle \phi(t_i) \phi(t_i) \rangle_c \end{aligned}$$

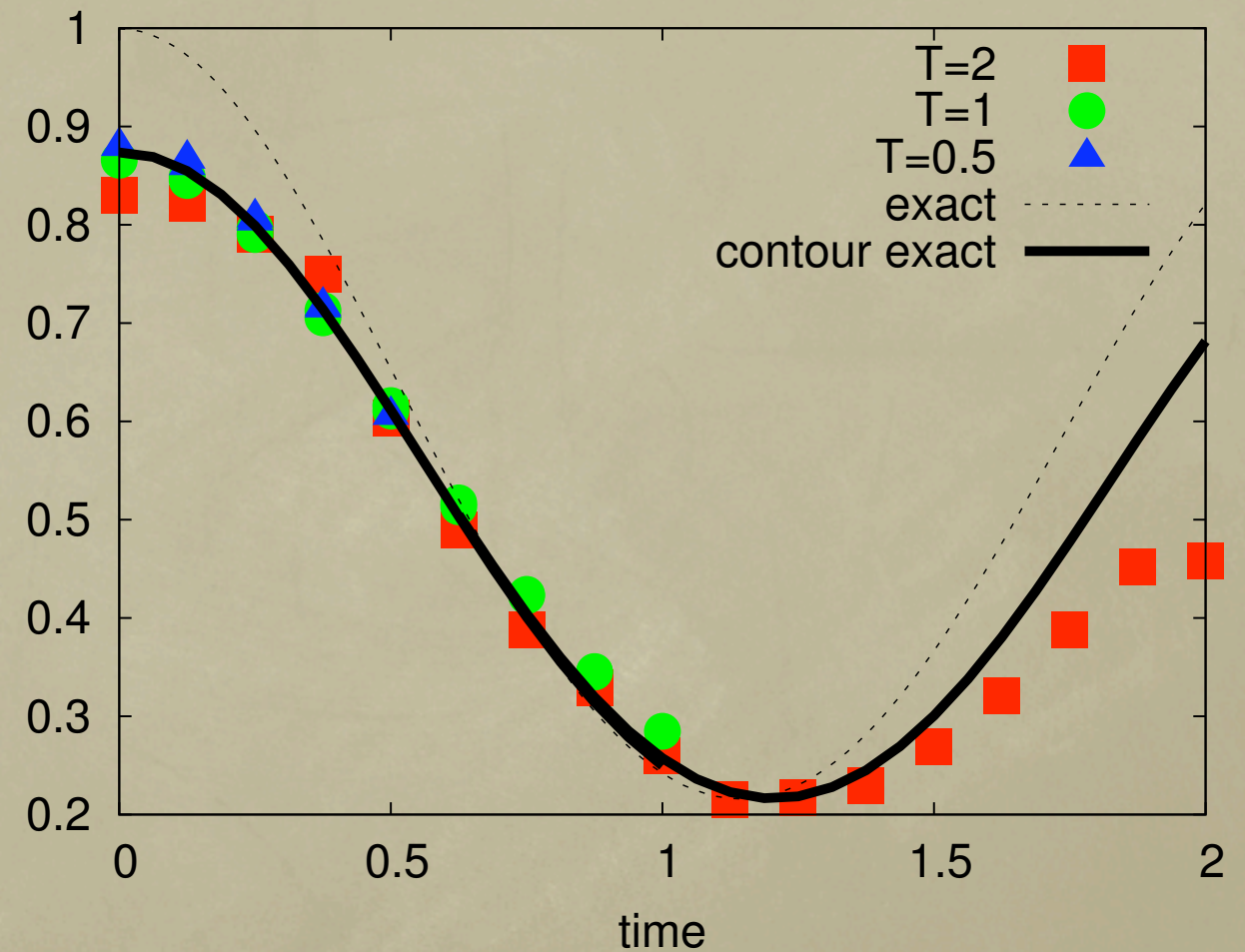


In practice:

$$\langle x(t) \rangle$$



$$\langle x(t)x(t) \rangle_c$$



*Reliable for short intervals*

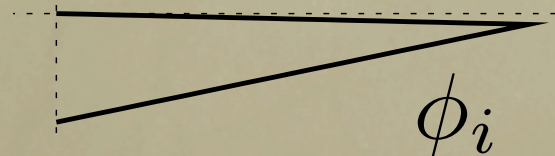


# Stochastic quantization:

real time simulations are indeed possible,

- Direct simulation on short intervals so far  
*(longer at low temperature or weak coupling)*

- For statistical estimate of real time correlators:  
more relevant information than Euclidean  
simulation may provide.



- Early nonequilibrium behaviour (instabilities)